2 光子励起によるX 線散乱



内殻分光

• 共鳴過程

物理現象が強調されて観測される 通常観測されない現象が観測される 多体効果などが観測される

• 軟X線

高輝度必須(要Undulator)

→ 新リング

• 硬X線

Bending 可

- → 開発研究多機能ビームライン
- 内殻エネルギー

遷移金属 *K*端: 4500~9000 eV 遷移金属 *L*端: 400~1000 eV 軽元素*K*端: 50~ 700 eV

X線発光·散乱

X-ray Emission Spectroscopy





 $TiO_2 (3d^0)$

 $L\alpha_{1,2}$

 \leftarrow $L\beta_1$

 $L\alpha_{1,2}$

 $\mathcal{L}\beta_1$

 $Ti_2O_3(3d^1)$





硬X線ラマン散乱





Y. Tezuka, et al., Phys. Rev B104, 235148 (2022)



2光子励起

- •1内殻同時励起 非線形ラマン散乱
- 2内殻励起 電子相関 原子間相関
- 誘導輻射
 素励起の強調
 ポンププローブ

2光子励起(非線形ラマン散乱)

Hyper-Raman Scattering



」、「「」、」、「」、「」」、「」」 通常のラマン散乱で観測できない素励起の観測

cf. Raman Scattering

△ℓ=0,2:単極子、4重極子

要

- 光子密度
- •時間コヒーレンス

cf. KEK 放射光 Conceptual Design Report (2016) 非線形非弾性X線散乱 (2光子吸収による非弾性散乱) エキゾティック相転移

水晶: α-β相転移 (573°C)

α水晶:三方晶系(trigonal System)

*β*水晶:六方晶系(hexagonal System)

D ₃	E	2C ₃	3C ₂	基底 変位、回転	基底 分極率
A ₁	1	1	1		$\alpha_{xx} + \alpha_{yy}, \alpha_{zz}$ β_{xxx}, β_{yyx}
A ₂	1	1	-1	T _z , R _z	$\beta_{xxy}, \beta_{xxz}, \beta_{yyy}, \beta_{yyz}, \beta_{zzz}$
E	2	-1	0	$(T_{x},T_{y}), (R_{x},R_{y})$	$\begin{array}{c} (\alpha_{xx} - \alpha_{yy}, \alpha_{xy}), (\alpha_{yz}, \alpha_{zx}) \\ \beta_{xxx}, \beta_{xxy}, \beta_{xxz}, \beta_{yyx}, \beta_{yyy}, \beta_{yyz}, \\ \beta_{zzx}, \beta_{zzy}, \beta_{zzz}, \beta_{xyz} \end{array}$

SHG禁制

D ₆	E	2C ₆	2C ₃	C ₂	3C ₂ '	3C ₂ ''	基底 変位、回転	基底 分極率
A ₁	1	1	1	1	1	1		$\alpha_{xx}\text{+}\alpha_{yy}\text{,}\alpha_{zz}$
A ₂	1	1	1	1	-1	-1	T_z , R_z	eta_{xxz} , eta_{yyz} , eta_{zzz}
B ₁	1	-1	1	-1	1	-1		$oldsymbol{eta}_{xxx}$, $oldsymbol{eta}_{yyx}$
B ₂	1	-1	1	-1	-1	1		$m{eta}_{xxy}$, $m{eta}_{yyy}$
E1	2	1	-1	-2	0	0	(T _x ,T _y), (R _x ,R _y)	$(\alpha_{yz}, \alpha_{zx}) \\ \beta_{xxx}, \beta_{xxy}, \beta_{yyx}, \\ \beta_{yyy}, \beta_{zzx}, \beta_{zzy}, \beta_{zzx}$
E ₂	2	-1	-1	-2	0	0		$(\alpha_{xx} - \alpha_{yy}, \alpha_{zy})$ $\beta_{xxz}, \beta_{yyz}, \beta_{xyz}$





X線ラマン散乱(2内殻励起)



より高次の多体効果

Ni, Cu *etc* : $3d^9$

Coherent X-ray Raman Spectroscopy

S.Tanaka and S. Mukamel, PRL 89, 043001(2002)

Coherent X-Ray Raman Spectroscopy: A Nonlinear Local Probe for Electronic Excitations

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Nonlinear x-ray four-wave mixing experiments are becoming feasible due to rapid advances in high harmonic generation and synchrotron radiation coherent x-ray sources. By tuning the difference of two x-ray frequencies across the valence excitations, it is possible to probe the entire manifold of molecular electronic excitations. We show that the wave vector and frequency profiles of this x-ray analogue of coherent Raman spectroscopy provide an excellent real-space probe that carries most valuable structural and dynamical information, not available from spontaneous Raman techniques.



FIG. 1. (a) One-body picture of CXRS. The open and hatched blocks are unoccupied conduction and unoccupied valence states, respectively. (b) One-dimensional molecular chain model system. The strengths of the interactions are written (in eV) and the site energies are given in parentheses. (c) Level scheme for CXRS.



FIG. 2. The calculated optical absorption (a) and x-ray absorption spectra (b).







FIG. 4. Wave vector dependence of CXRS when $\omega_1 = 104.05 \text{ eV}$ and $\omega_s = 203.77 \text{ eV}$. The grating wave vectors are $q = 2\pi u/|\mathbf{R}_a - \mathbf{R}_b|$, for various values of u as indicated.



FIG. 5. Wave vector dependence of CXRS for the exciton peaks of Fig. 4.

軟X線における2ビーム利用

- 高輝度な2-ビーム
- Coherence



Undulator $\times 2$

蛍光X線



軟X線の増幅

内殻吸収

